

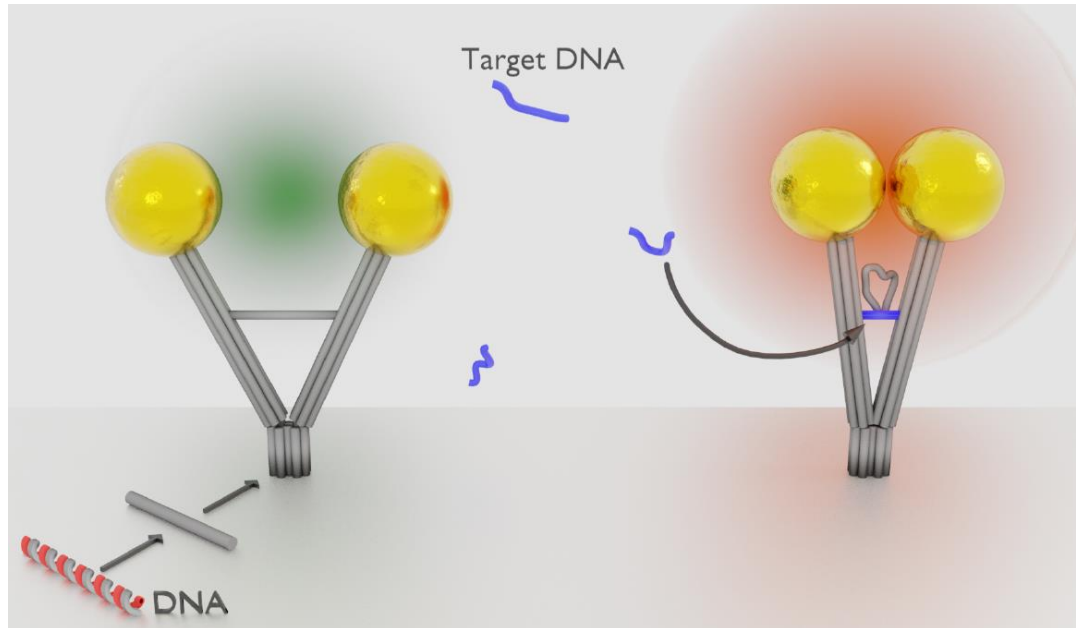
# Structural DNA nanotechnology as a playground for plasmonics

Sébastien Bidault

Institut Langevin, ESPCI Paris, CNRS, PSL Research University, Paris, France  
[http://www.institut-langevin.espci.fr/optical\\_antennas](http://www.institut-langevin.espci.fr/optical_antennas)



# Why structural DNA nanotechnology?



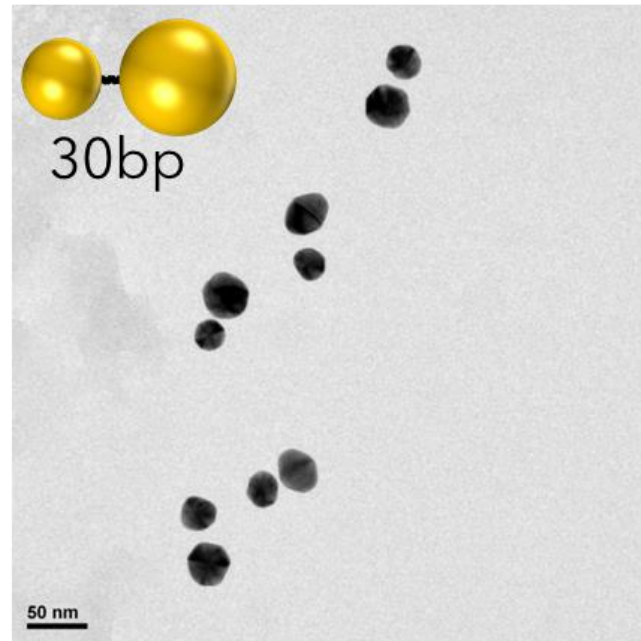
- Producing dynamic nanostructures whose morphology and optical properties can be actively and specifically modulated

- Introducing a controlled number of quantum emitters in the hot-spot of a plasmonic resonator

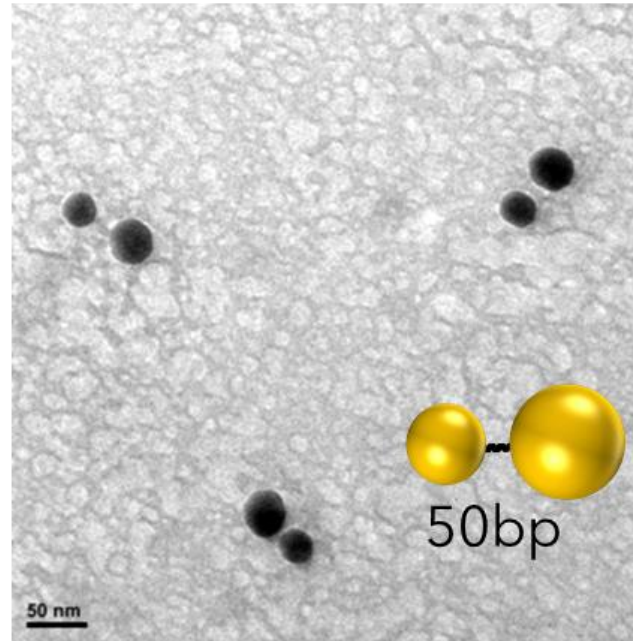


# DNA linked gold nanoparticle dimers

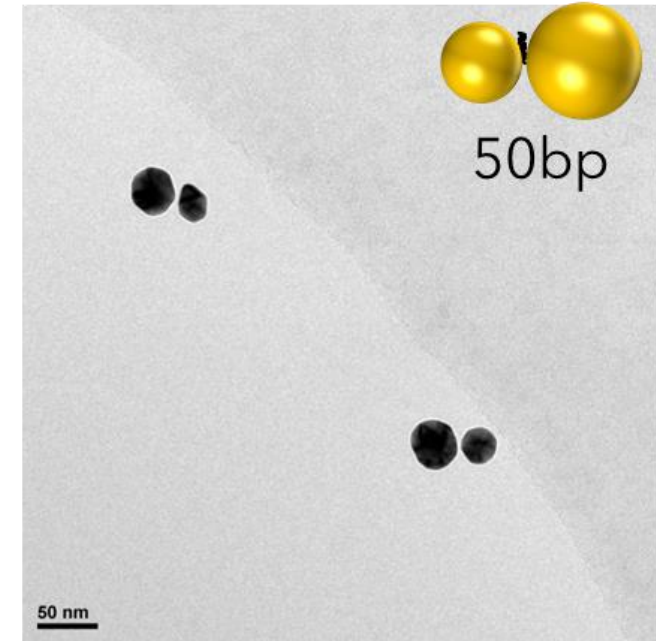
40 nm / 30 nm Au dimers:



$d=13\pm2$  nm



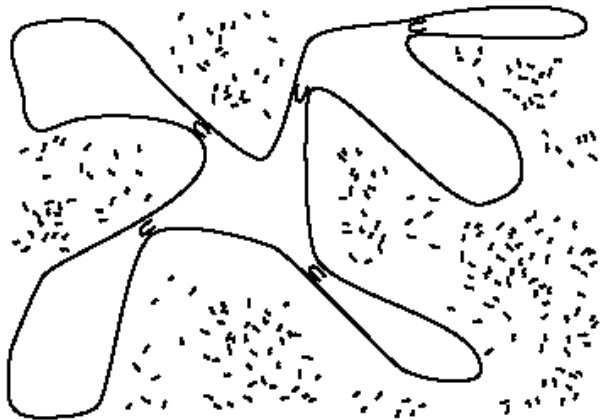
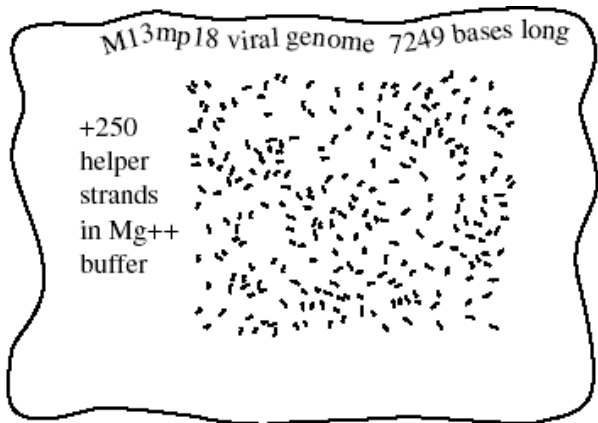
$d=17.5\pm3$  nm



$d=7\pm1$  nm

cryo-EM imaging (to avoid drying effects)

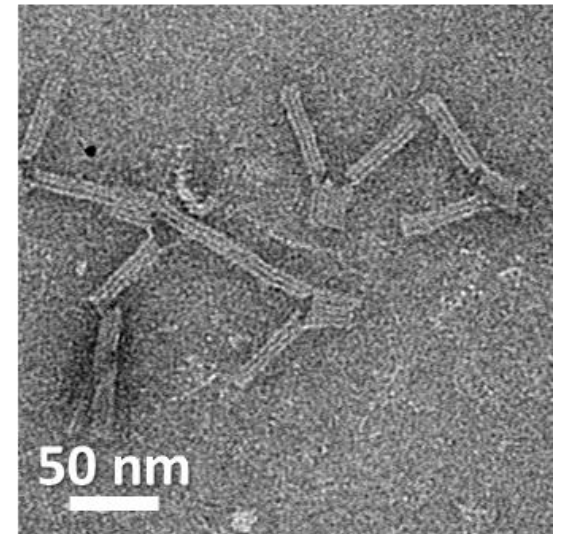
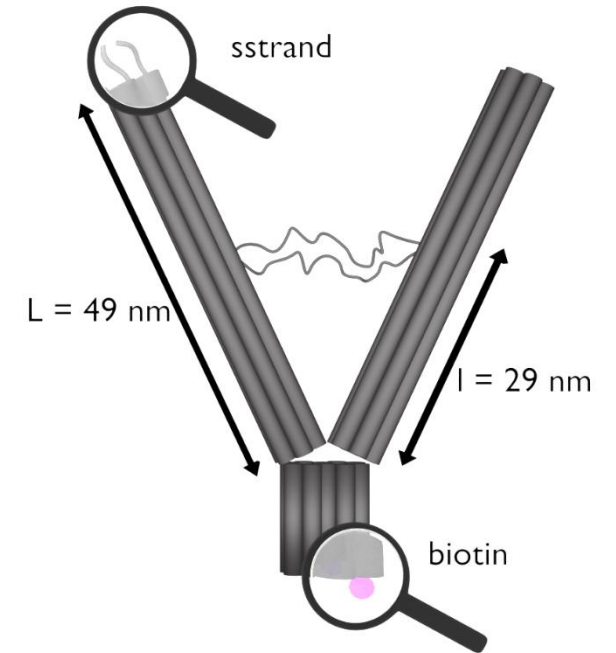
# DNA origami templates



P.K. Rothemund, Nature 2006



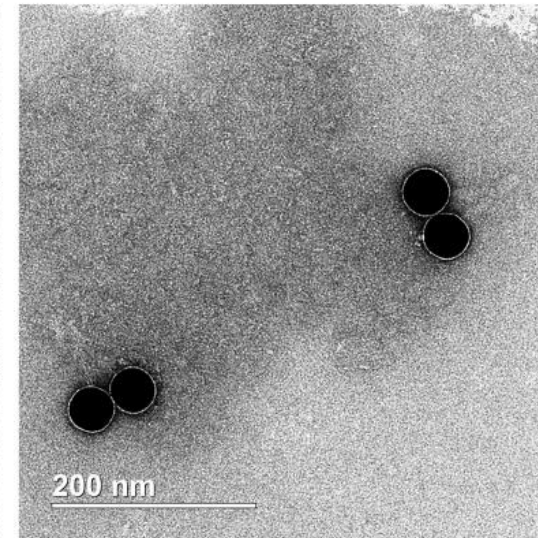
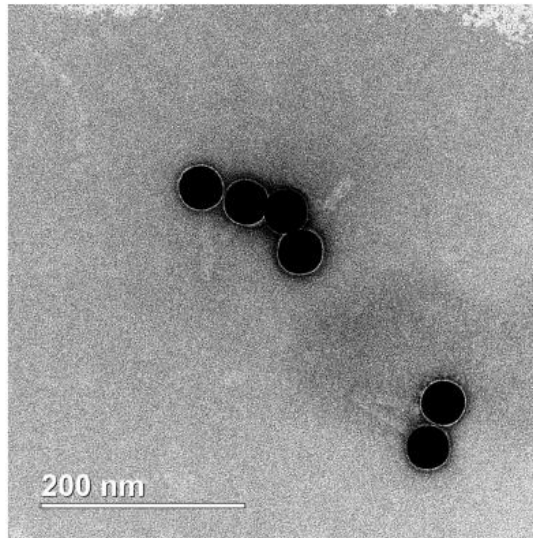
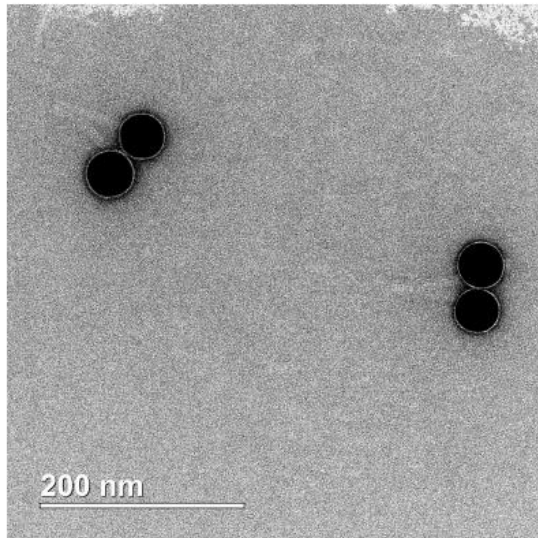
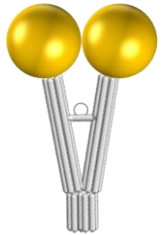
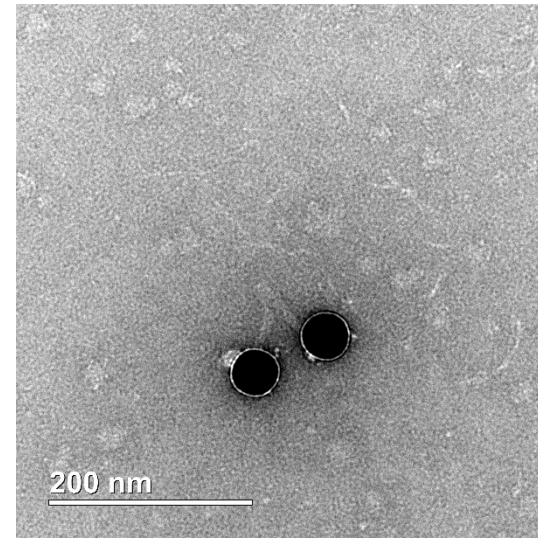
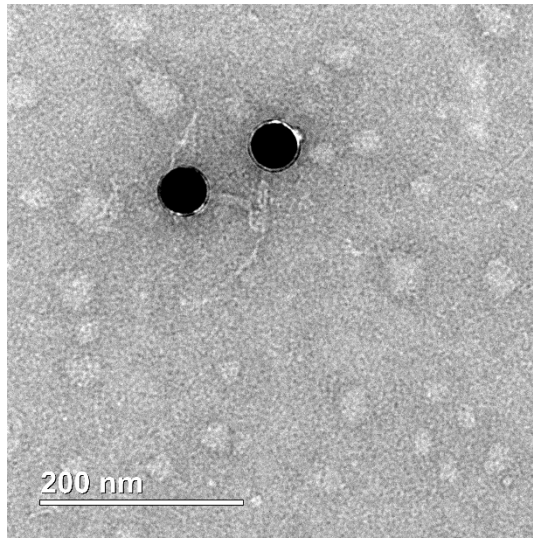
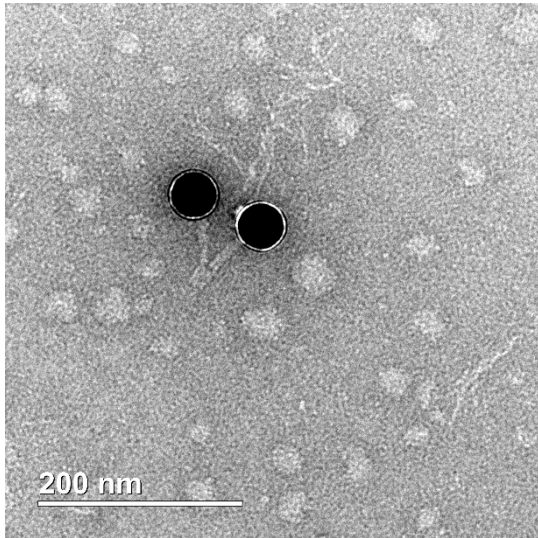
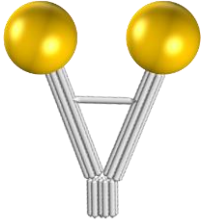
S. M. Douglas et al, Nature 2009 & H. Dietz et al, Science 2009



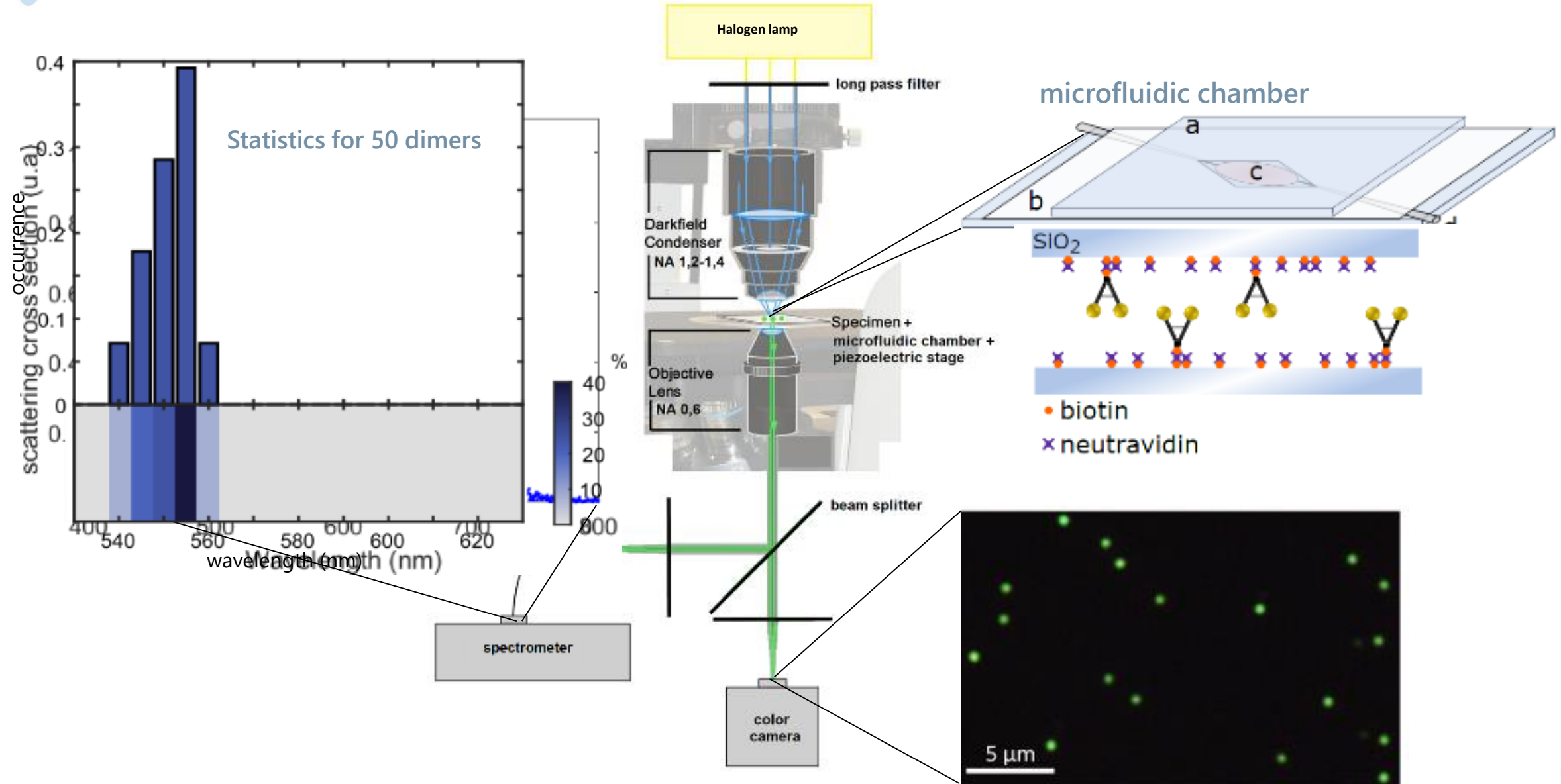




# Gold nanoparticle dimers on DNA-origami templates



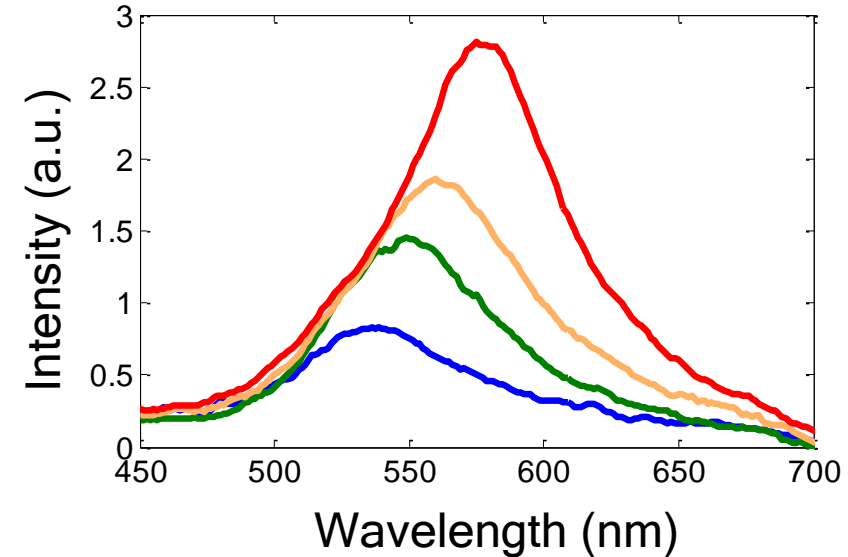
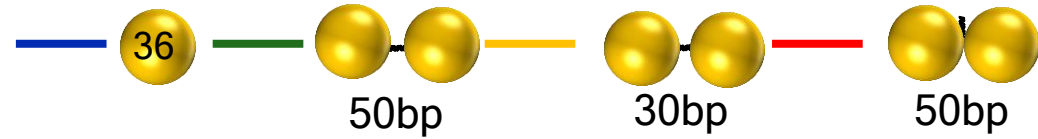
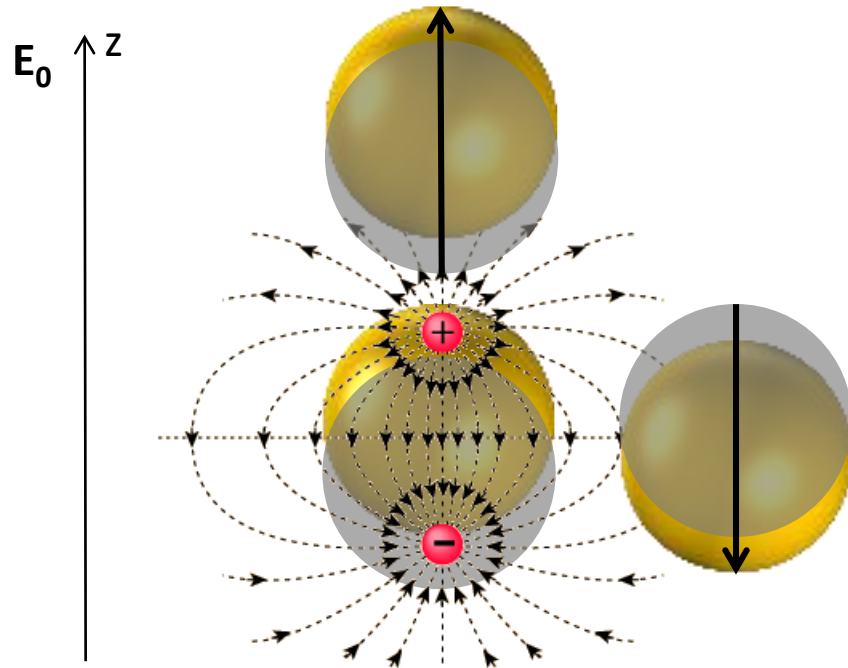
# Scattering spectroscopy







# Near-field plasmon coupling



Reducing the interparticle distance induces a wavelength redshift of the longitudinal plasmon mode of the dimers and increases the scattering cross-section of the antenna

Monitoring the longitudinal plasmon resonance allows a **nanoscale analysis of the interparticle distance**

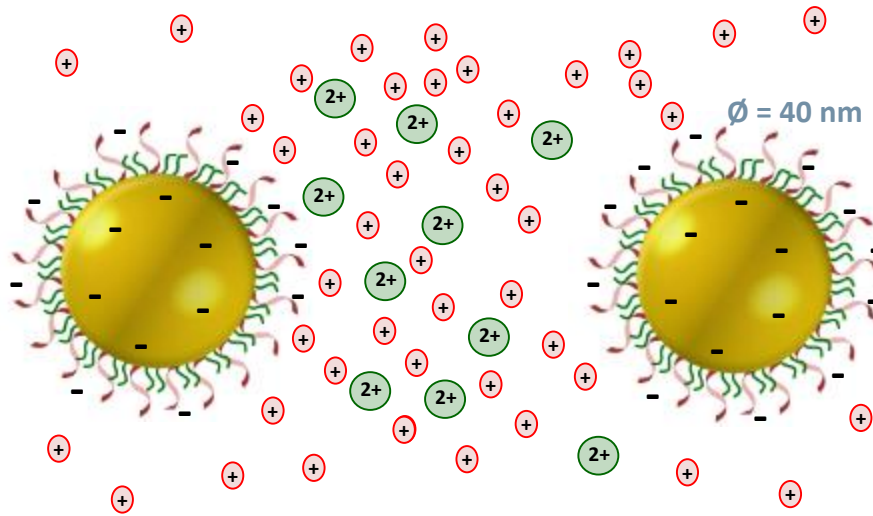
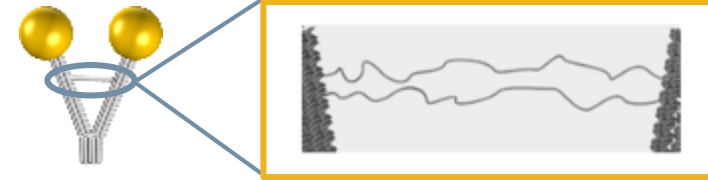


# Actively tuning the interparticle distance

Ionic strength [M]

$$I = \frac{1}{2} \sum_i c_i z_i^2$$

Single-stranded active site



2+  $Mg^{2+}$

+  $Na^{+}$



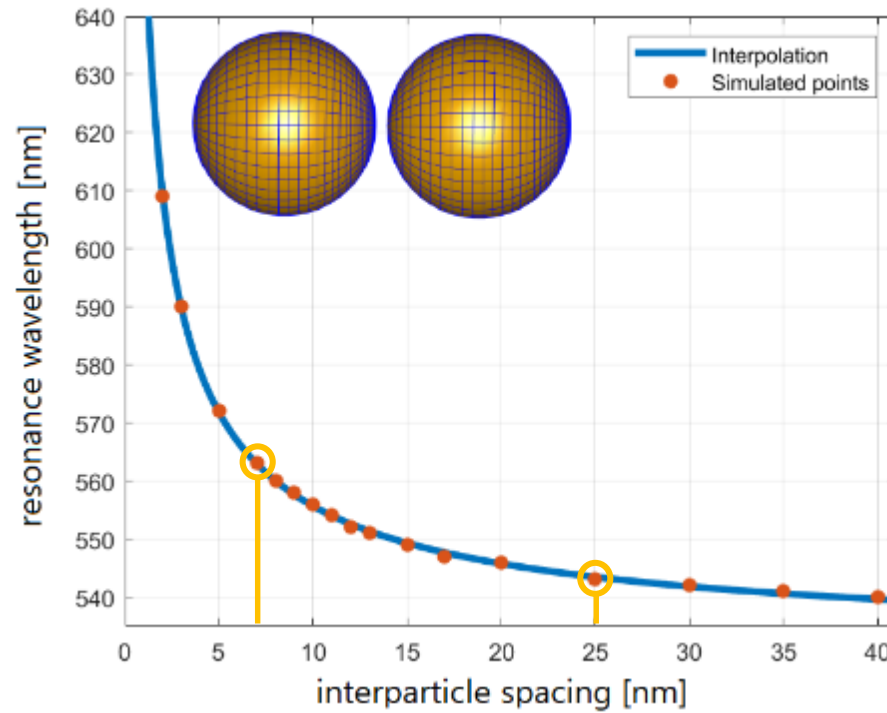


# Actively tuning the interparticle distance

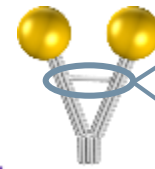
Ionic strength [M]

$$I = \frac{1}{2} \sum_i c_i z_i^2$$

calibration curve

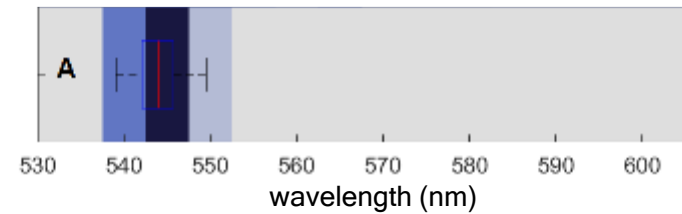


Single-stranded active site



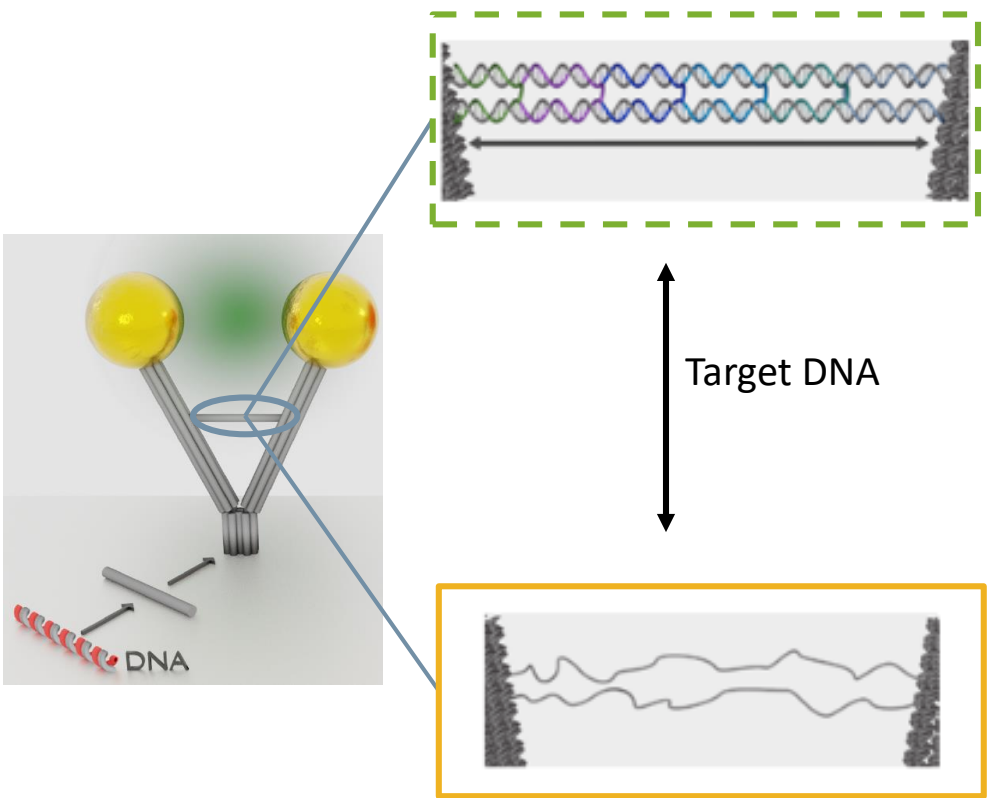
I  
(mM)

25  
1655  
~15h  
1655  
455  
25

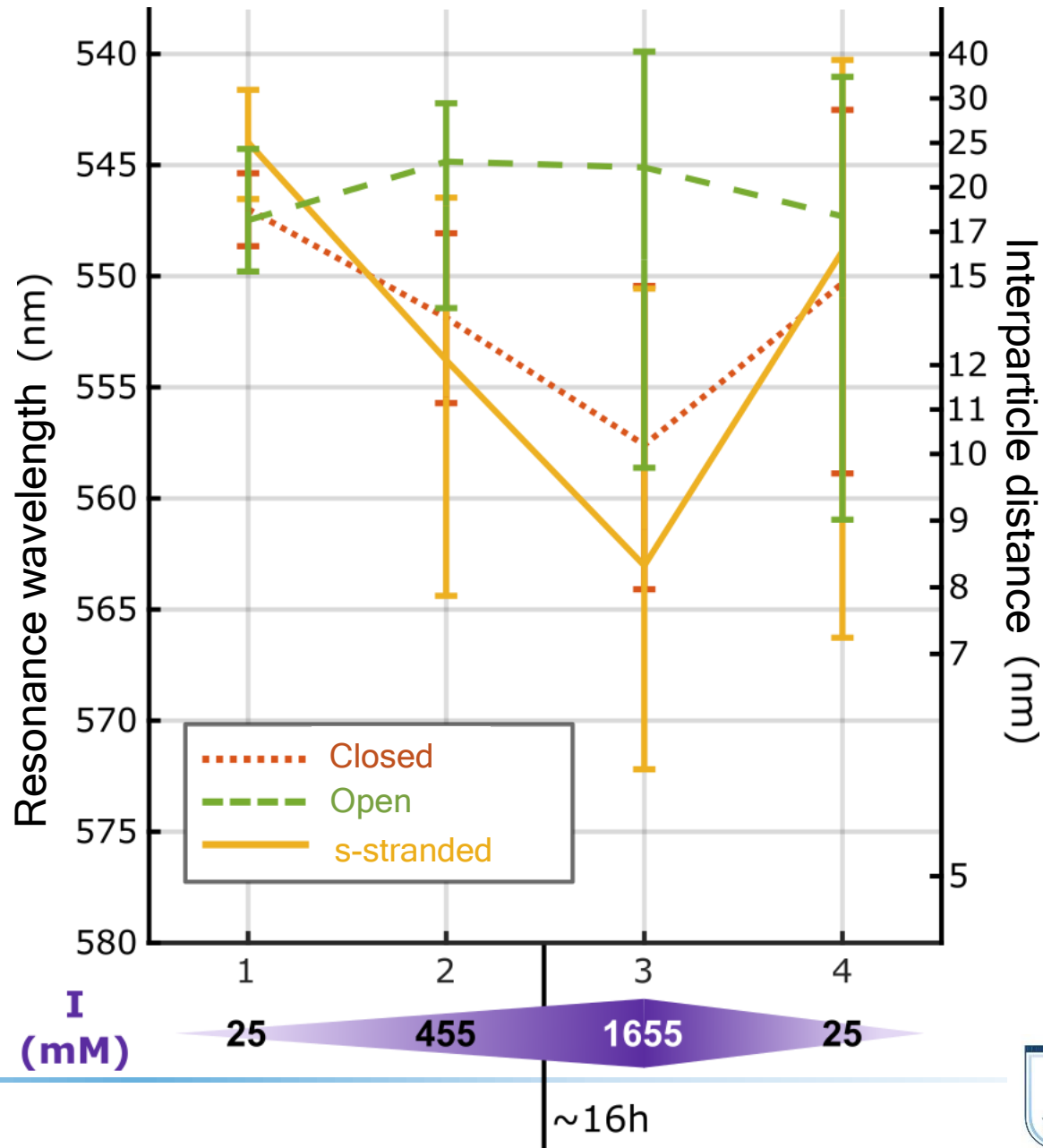




# Tuning the rigidity of the origami

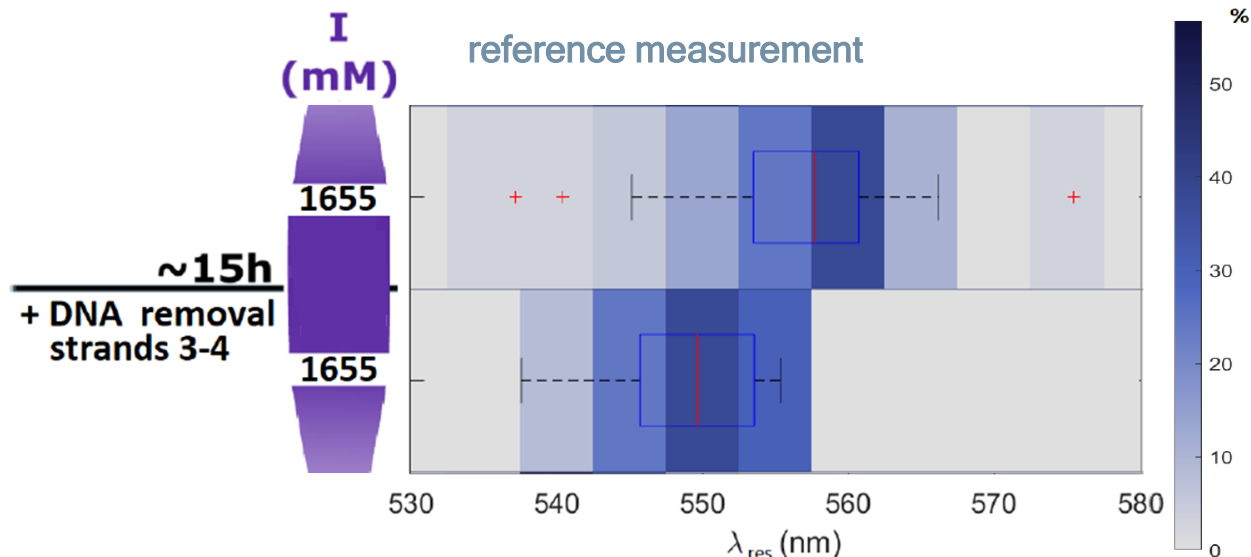
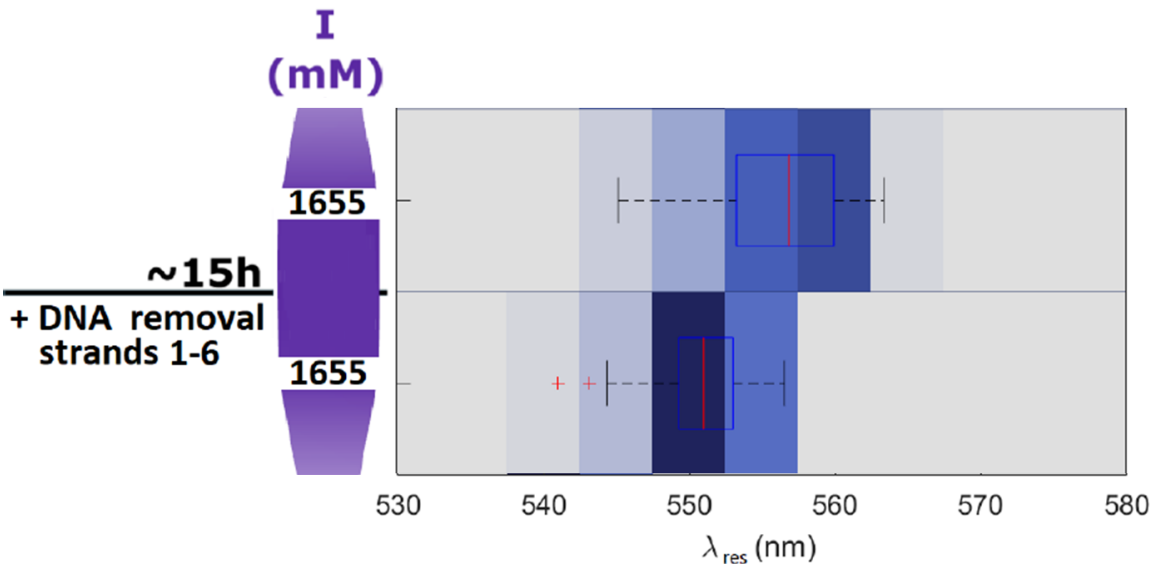
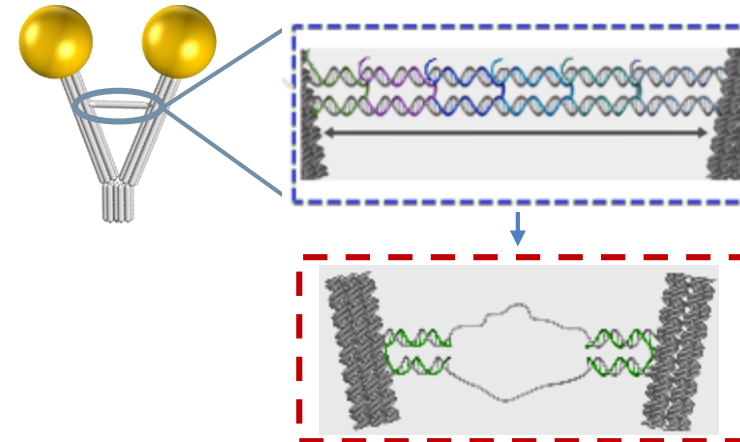
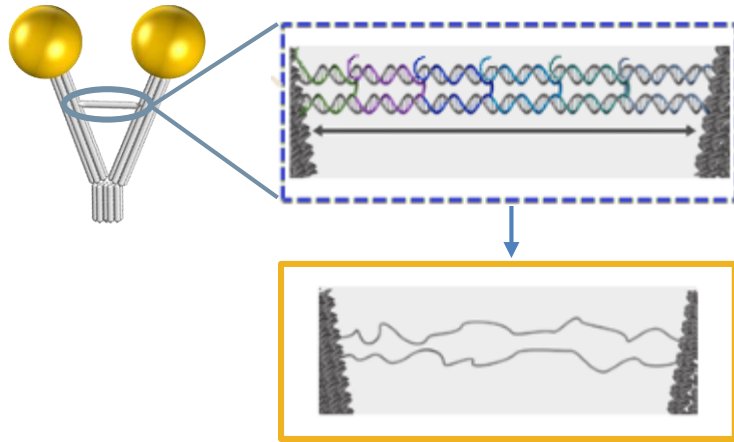


The geometry and rigidity of the DNA origami control the flexibility of the hybrid nanostructure





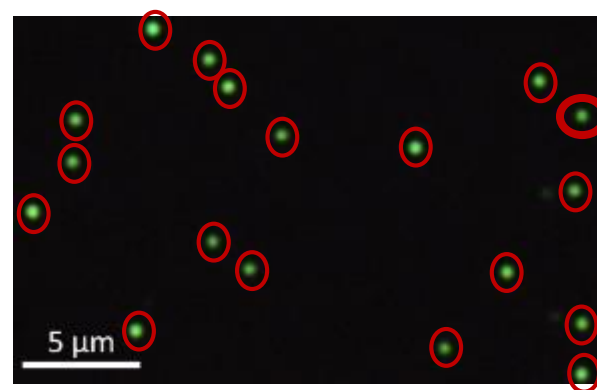
# Optical detection of a DNA strand displacement reaction



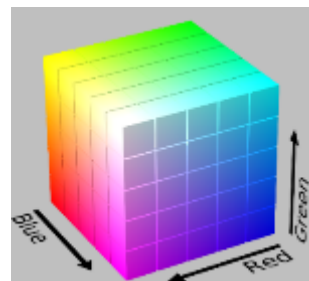
Strand displacement induces spectral redshift of  $\sim 7\text{nm}$



# Colorimetric detection of a DNA strand displacement reaction

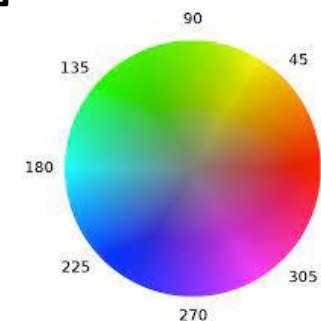
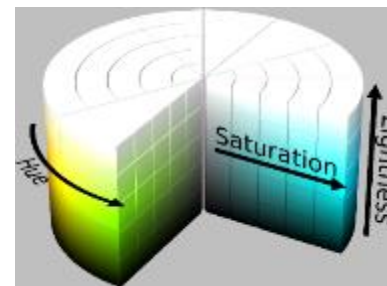


[R,G,B]



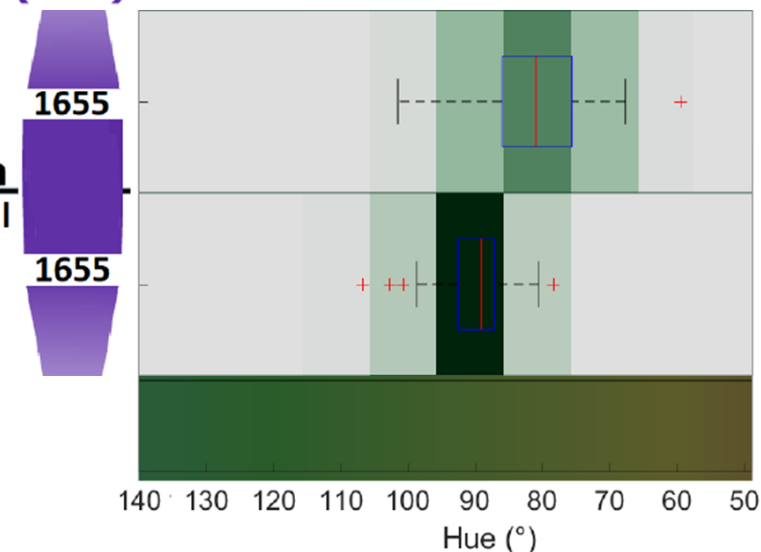
→

[H,S,L]



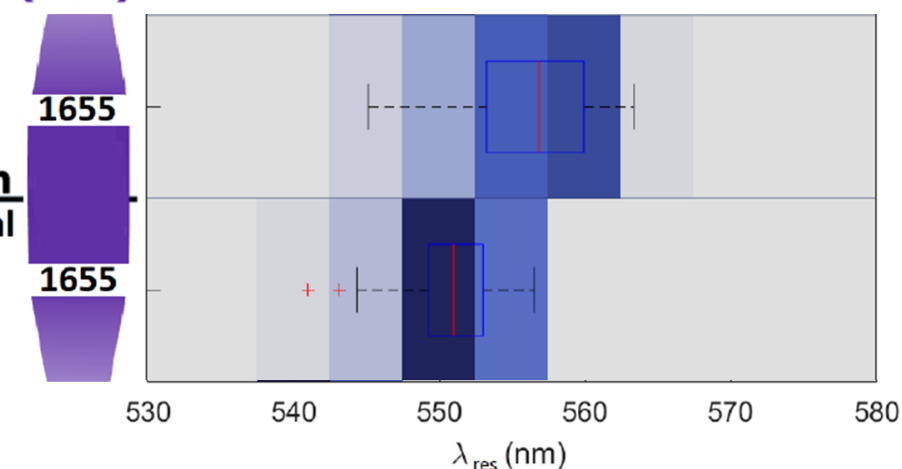
I  
(mM)

**~15h**  
+ DNA removal  
strands 1-6



I  
(mM)

**~15h**  
+ DNA removal  
strands 1-6

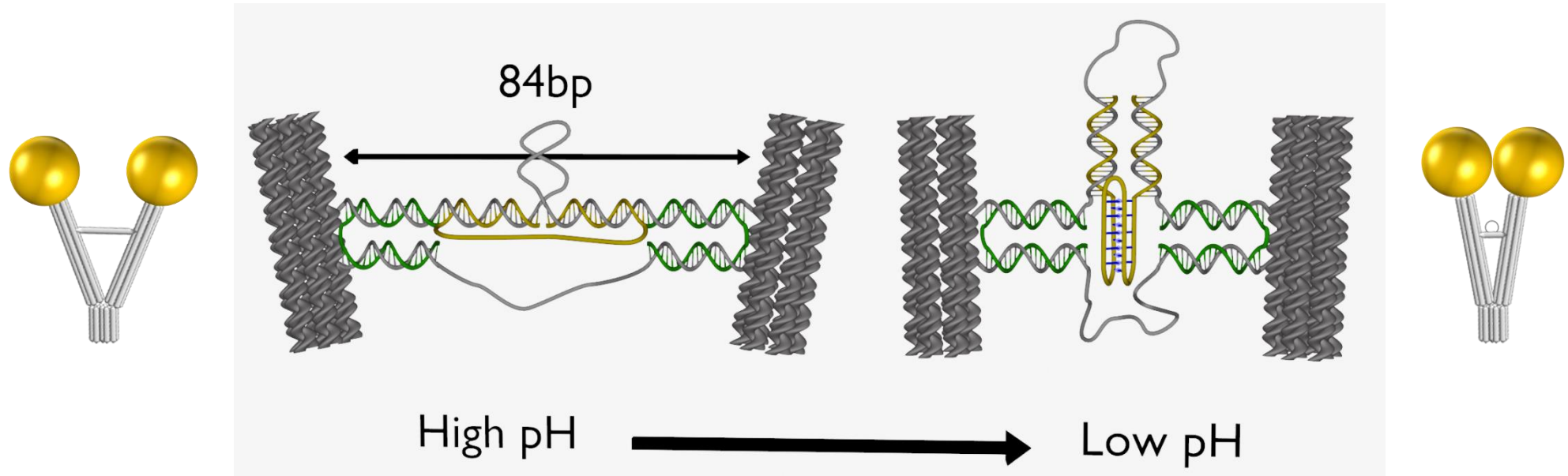


Colorimetric detection of strand displacement similar to data collected via spectrometer

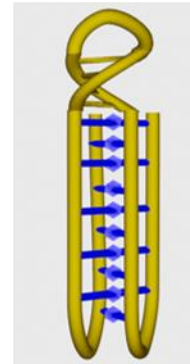
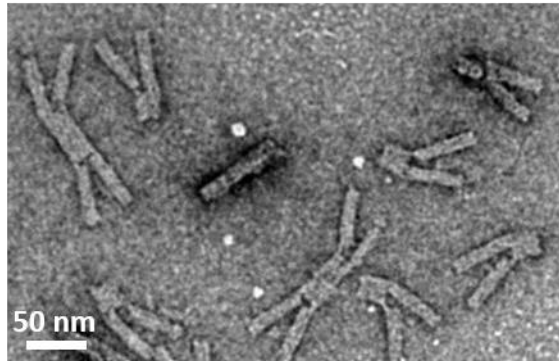




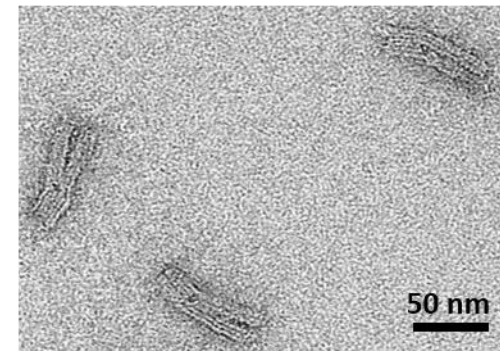
# pH-sensitive hybrid nanostructures



pH  
8.5



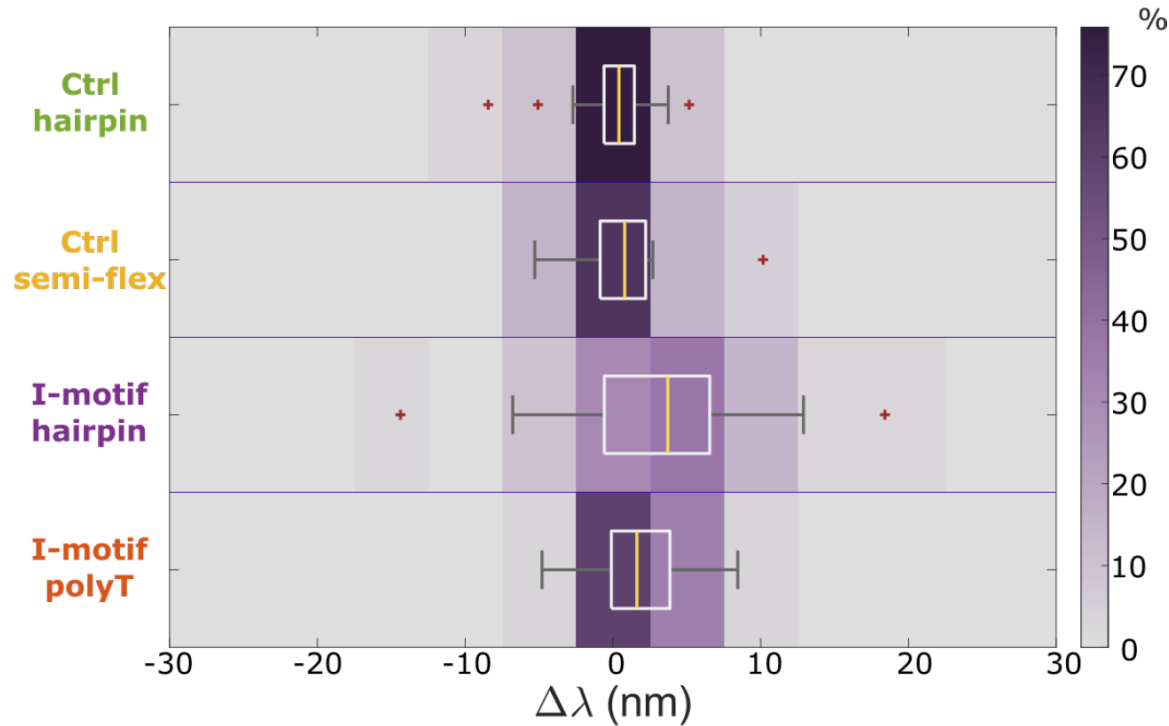
Hairpin:  
4 bp (1 AT, 3 GC)



pH  
5.6

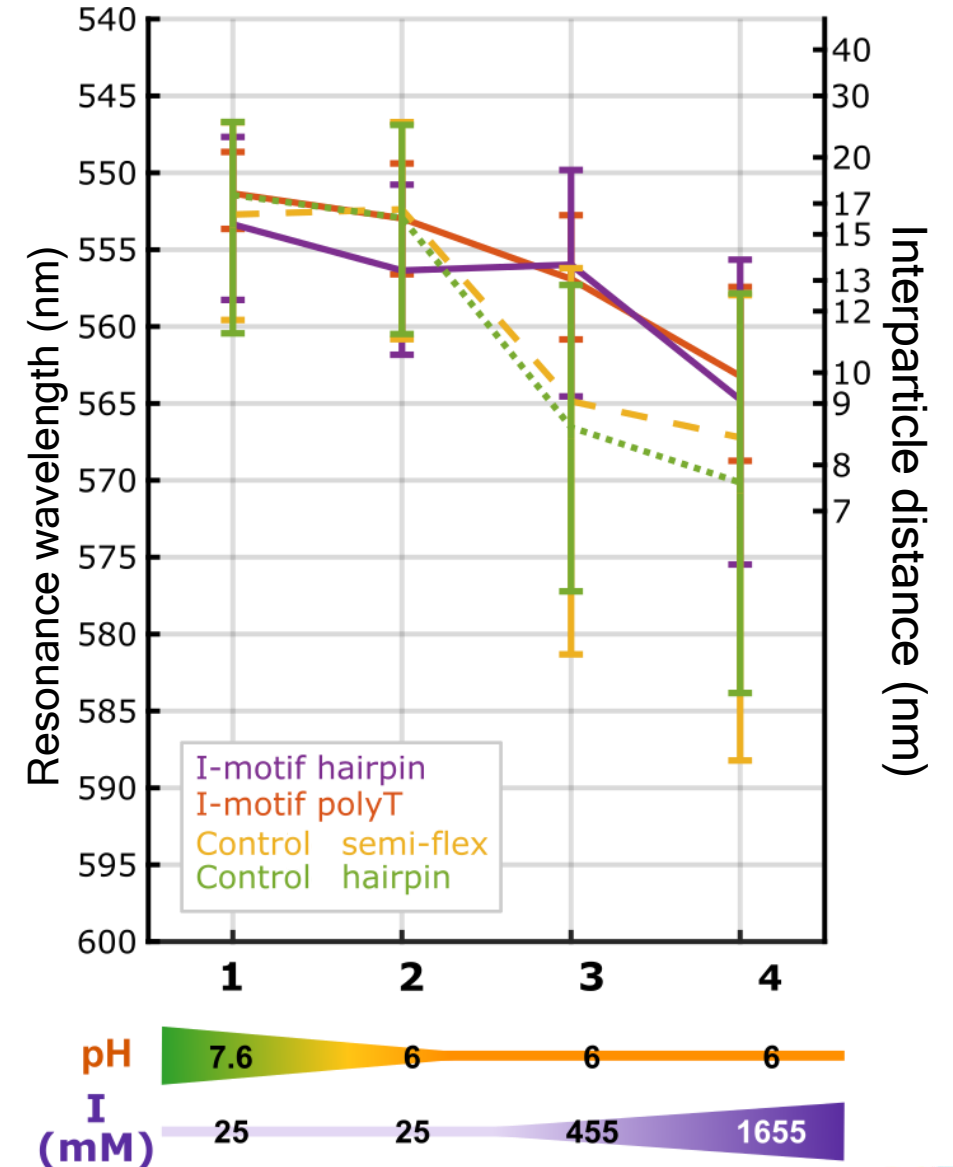


# pH-sensitive hybrid nanostructures



Small redshift due to DNA i-motif during the pH transition

Large modification of the rigidity of the hybrid nanostructure





# Why structural DNA nanotechnology?



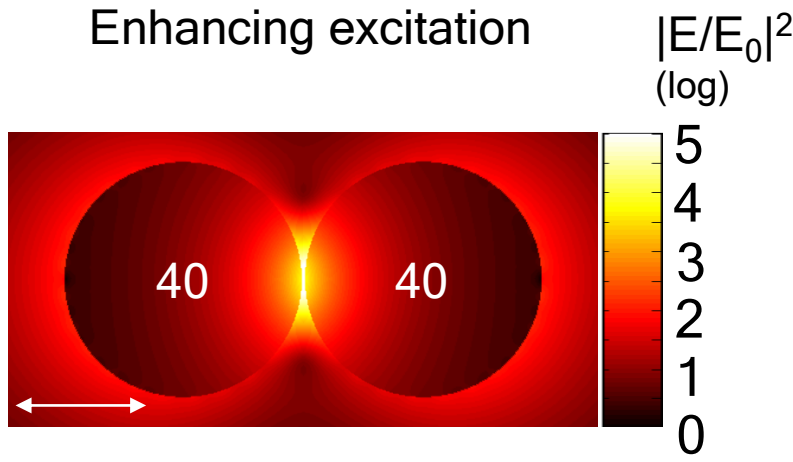
- Producing dynamic nanostructures whose morphology and optical properties can be actively and specifically modulated

- Introducing a controlled number of quantum emitters in the hot-spot of a plasmonic resonator



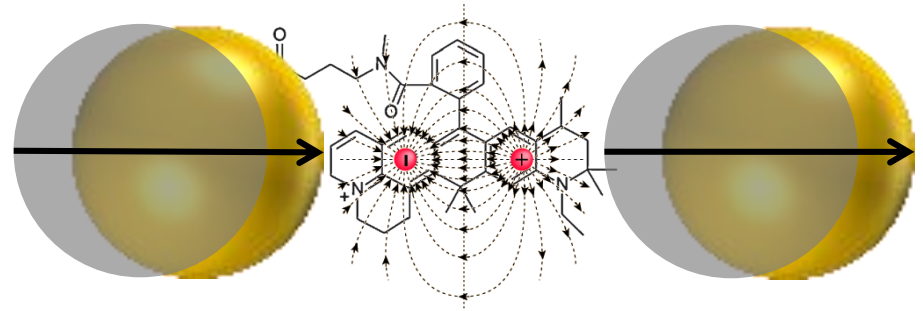
# Coupling broadband emitters and resonators

Enhancing excitation



$$P_{0 \rightarrow 1} = \sigma_{abs}(\omega_{exc}) \cdot |\mathbf{E}(\mathbf{r}) \cdot \mathbf{u}_p|^2$$

Modifying the local density of states



$$F_p = \frac{\rho_p(\mathbf{r}, \omega)}{\rho_0} = \frac{P}{P_0} = \frac{\Gamma}{\Gamma_0} \quad \begin{matrix} \text{(classical} \\ \text{dipole)} \end{matrix} \quad \begin{matrix} \text{(quantum weak} \\ \text{coupling regime)} \end{matrix}$$

Reciprocity in weak coupling  $\Rightarrow \frac{|\mathbf{u}_p \cdot \mathbf{E}|^2}{|\mathbf{u}_p \cdot \mathbf{E}_0|^2} \approx \frac{\Gamma_R}{\Gamma_{R0}}$

P. Bharadwaj et al, Adv. Opt. Photon. 2009

Enhancing the fluorescence signal

$$\frac{I_{fluo_{enhanced}}(\mathbf{r}, \mathbf{u}_p)}{I_{fluo_{ref}}} = \frac{|\mathbf{E}(\mathbf{r}) \cdot \mathbf{u}_p|^2}{|\mathbf{E}_0 \cdot \mathbf{u}_p|^2} \cdot \frac{\varphi(\mathbf{r}, \mathbf{u}_p)}{\varphi_0}$$





# Coupling regimes

Emitter ( $\omega_0, \gamma_0$ ) and resonator ( $\omega_p, \gamma_p$ ) as two coupled damped harmonic oscillators



Coupling rate  $g$ : complex eigenfrequencies of the coupled system are inferred by diagonalizing

$$\begin{pmatrix} \omega_0 - i(\gamma_0 - \Gamma_0)/2 & g \\ g & \omega_p - i\gamma_p/2 \end{pmatrix} \quad \text{with } g = \sqrt{F_p \Gamma_0 \gamma_p / 4}$$

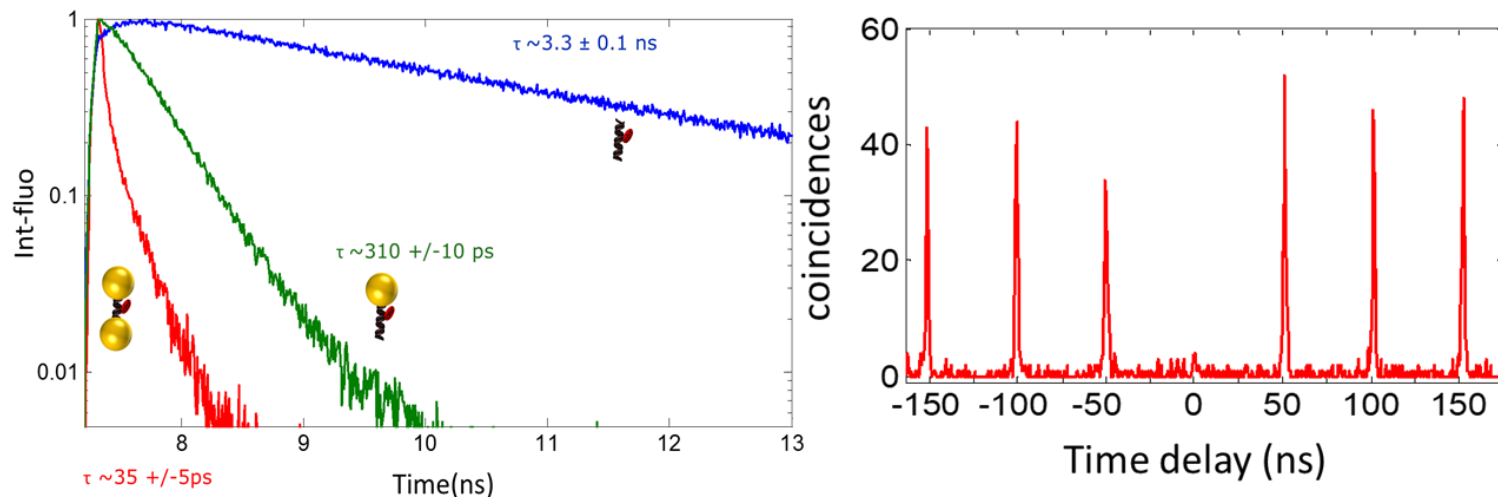
$$\text{Eigenmodes } \begin{cases} \omega_{\pm} = \frac{1}{2}(\omega_0 + \omega_p) \pm \sqrt{\left(\frac{\omega_0 - \omega_p}{2}\right)^2 + g^2} \\ \gamma_{\pm} = \frac{(\gamma_0 - \Gamma_0) + \gamma_p}{2} \end{cases}$$

Weak coupling regime if  $|\tilde{\omega}_0 - \tilde{\omega}_p|^2 \gg 4g^2$  then  $\omega_+ = \omega_0$  and  $\omega_- = \omega_p$

Strong coupling regime if  $|\tilde{\omega}_0 - \tilde{\omega}_p|^2 < 4g^2$

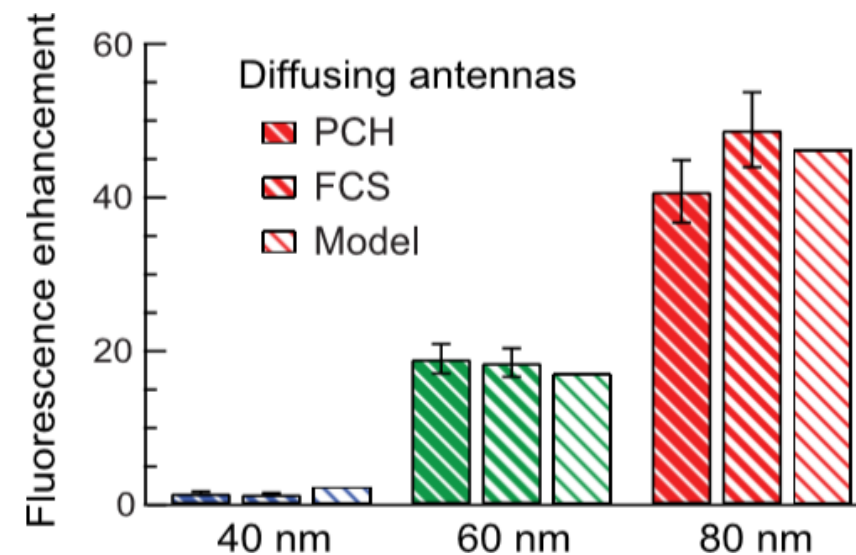
# Enhancing emission in the weak coupling regime

## Accelerating emission



M. P. Busson et al, Nat. Commun. 3, 962 (2012)

## Fluorescence enhancement



S. Bidault et al, ACS Nano , 10, 4806–4815 (2016)

$\phi \sim 15\%$      $\phi \sim 45\%$      $\phi \sim 70\%$

## High brightness and quantum yields

Fluorescence intensity up to  $\times 330$  ( $\times 44$  average)

Emission rates up to  $\times 760$  ( $\times 72$  average)

Typical antenna yields in 40 % - 70 % range

# Coupling regimes

Emitter ( $\omega_0, \gamma_0$ ) and resonator ( $\omega_p, \gamma_p$ ) as two coupled damped harmonic oscillators



Coupling rate  $g$ : complex eigenfrequencies of the coupled system are inferred by diagonalizing

$$\begin{pmatrix} \omega_0 - i(\gamma_0 - \Gamma_0)/2 & g \\ g & \omega_p - i\gamma_p/2 \end{pmatrix} \quad \text{with } g = \sqrt{F_p \Gamma_0 \gamma_p / 4}$$

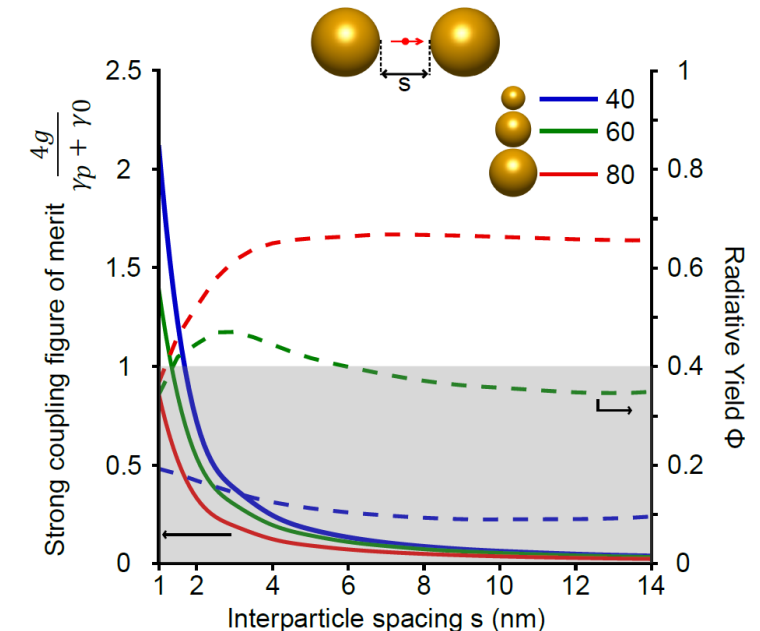
$$\text{Eigenmodes} \begin{cases} \omega_{\pm} = \frac{1}{2}(\omega_0 + \omega_p) \pm \sqrt{\left(\frac{\omega_0 - \omega_p}{2}\right)^2 + g^2} \\ \gamma_{\pm} = \frac{(\gamma_0 - \Gamma_0) + \gamma_p}{2} \end{cases}$$

Weak coupling regime if  $|\tilde{\omega}_0 - \tilde{\omega}_p|^2 \gg 4g^2$  then  $\omega_+ = \omega_0$  and  $\omega_- = \omega_p$

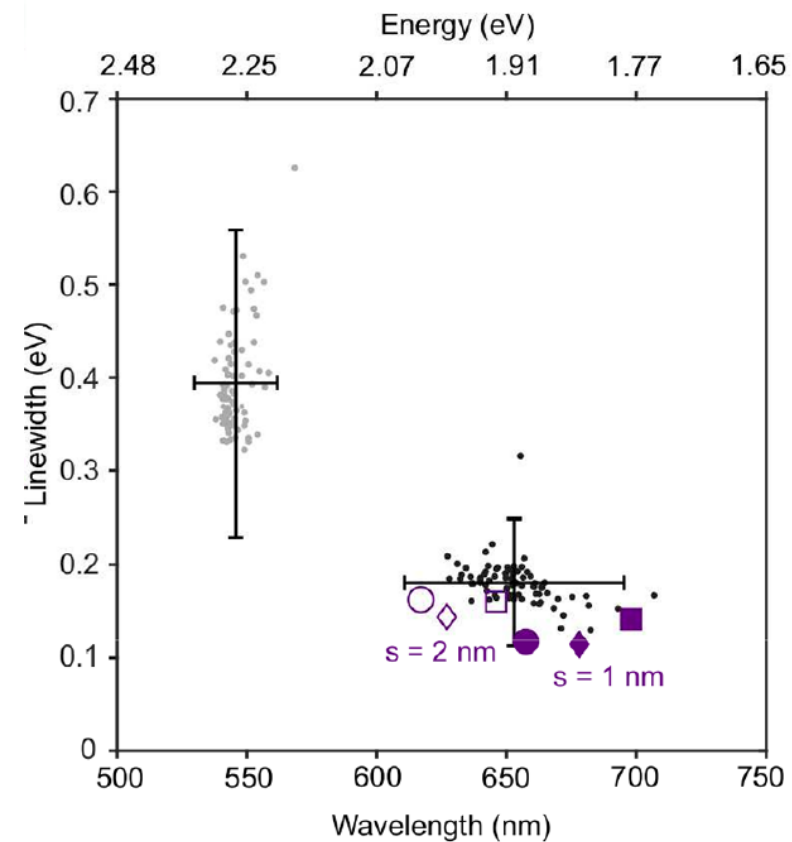
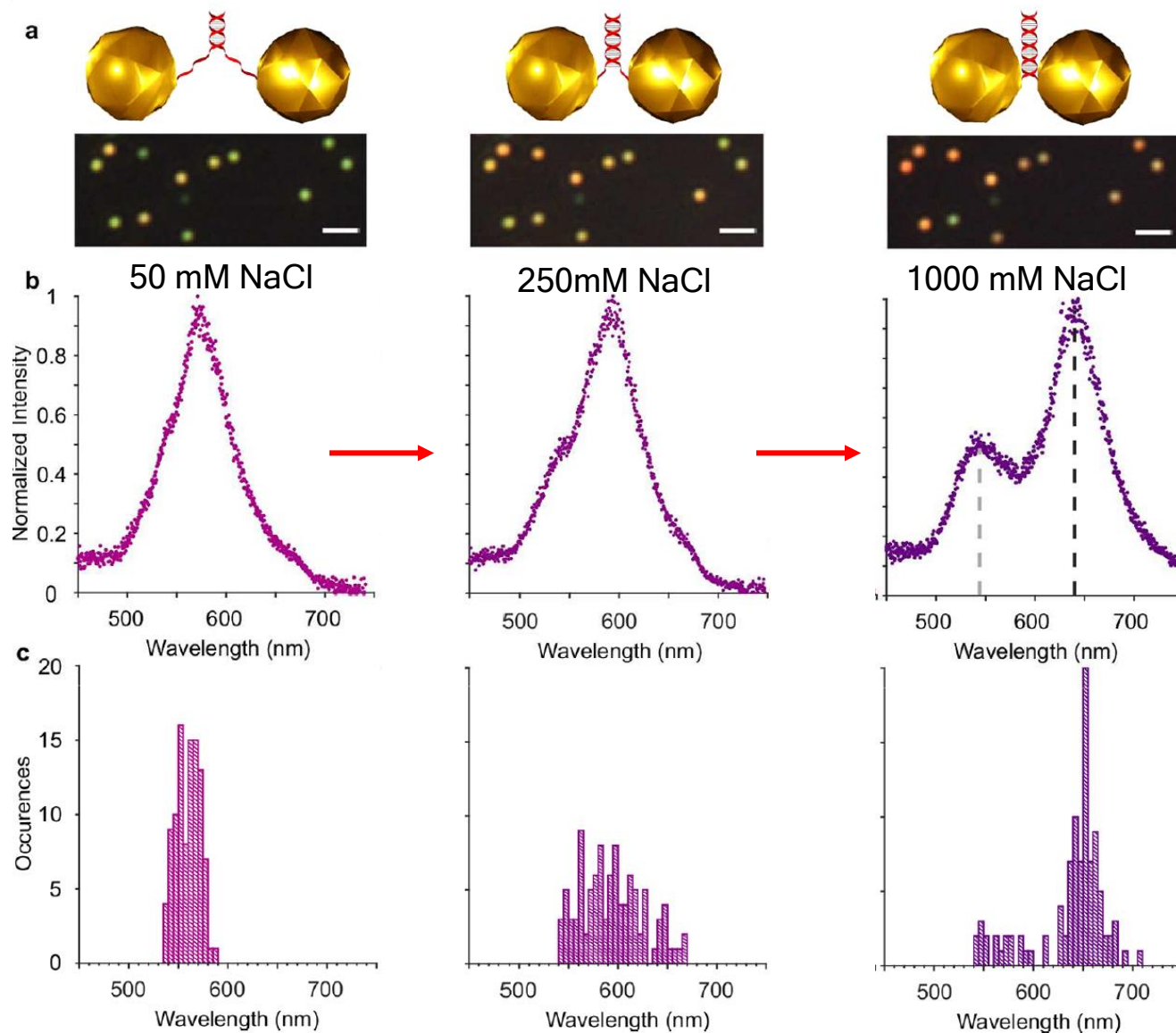
Strong coupling regime if  $|\tilde{\omega}_0 - \tilde{\omega}_p|^2 < 4g^2$

Two strongly coupled hybrid resonances are observed if the coupling rate is larger than all decay rates in the coupled system:  $2g > \gamma_p/2 + \gamma_0/2$

and with  $N$  identical emitters  $g_N = \sqrt{N}g$

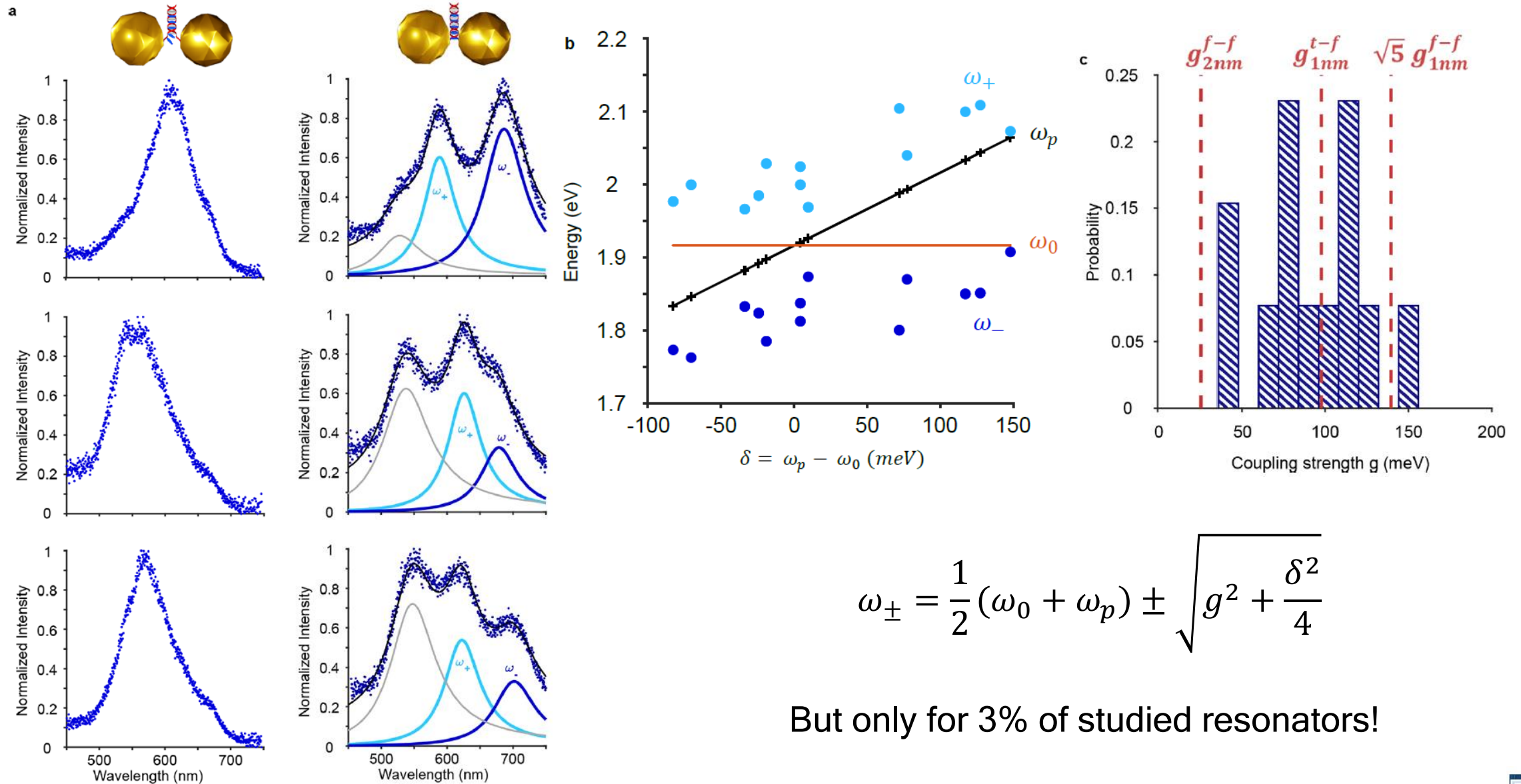


# Reaching sub-2nm gaps reproducibly

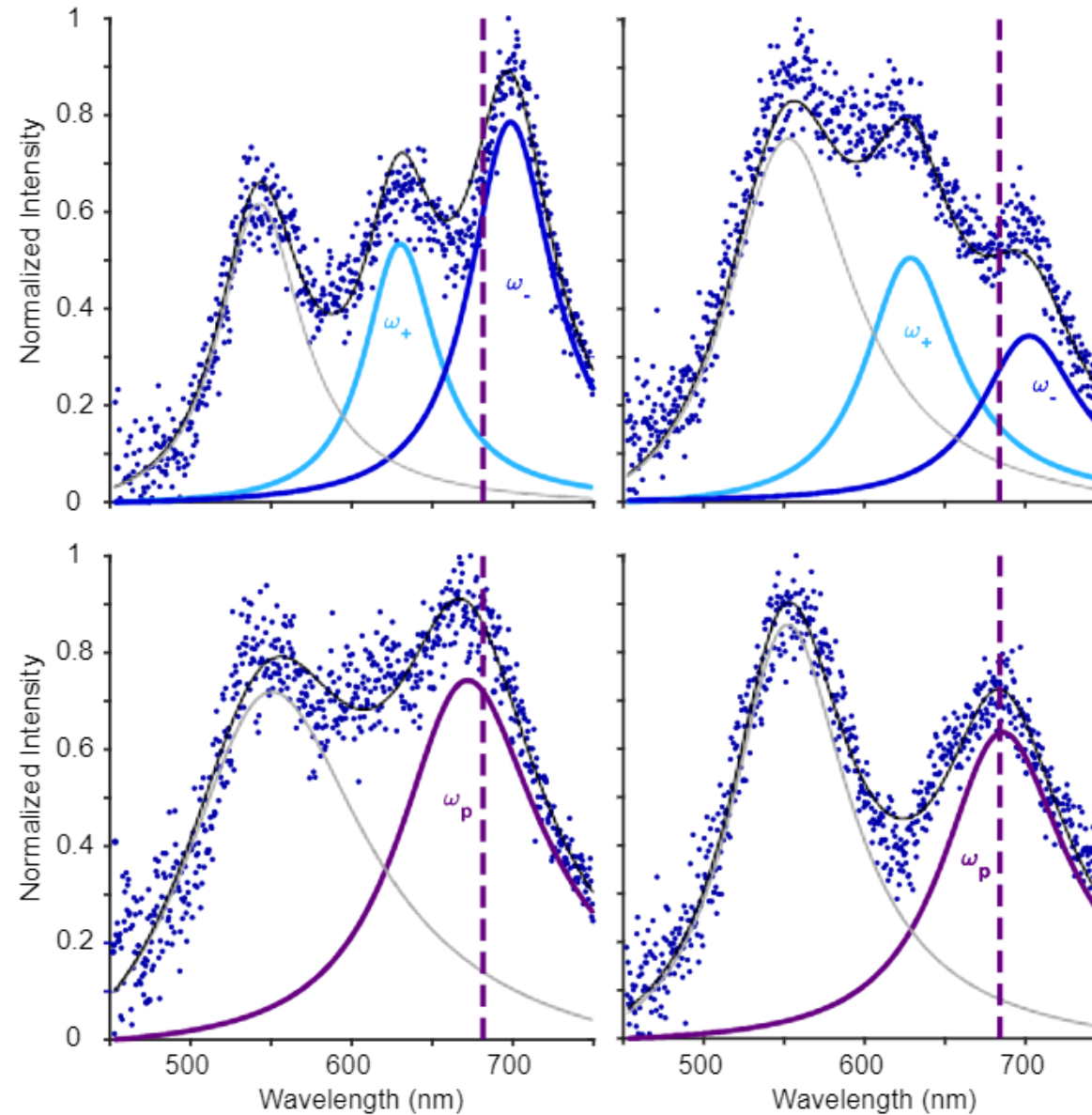
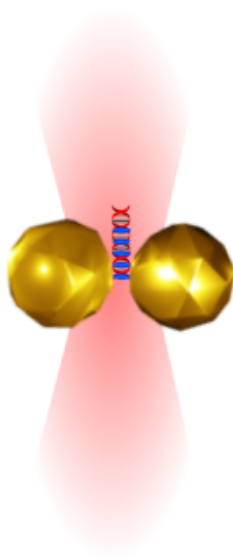




# Visible strong-coupling with 5 ATTO647N molecules

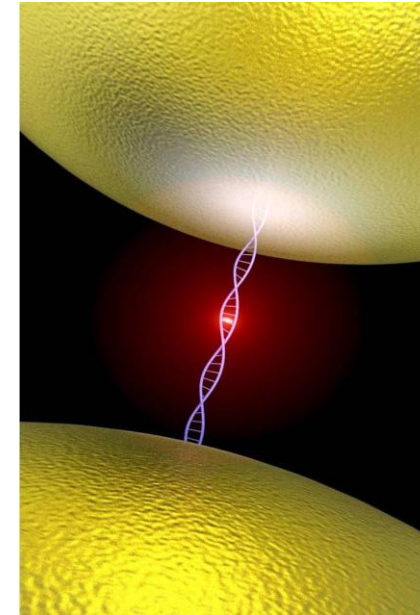
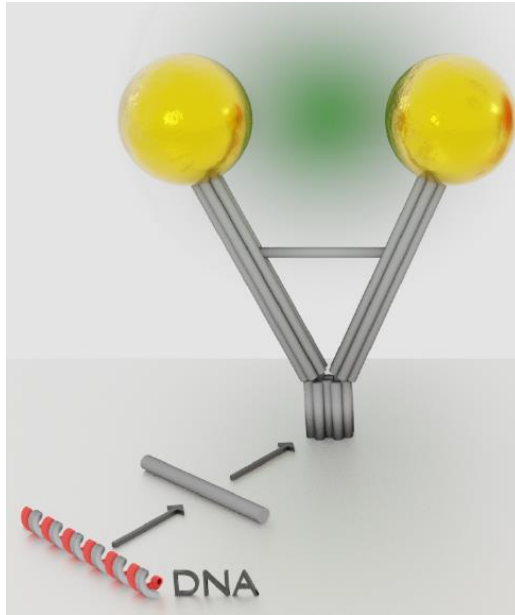


# Photoinduced disappearance of hybrid modes





# Conclusions



- DNA-based nanofabrication techniques allow the development of plasmonic resonators with an excellent control over their chemical environment
- An active control of the interparticle distance is possible by tuning interparticle interactions or the shape of the DNA scaffold
- DNA-templated dimers enhance both the fluorescence intensity and decay rate of single molecules by nearly 2 orders of magnitude (average) and up to 3 (max)
- It is possible to reach a strong-coupling regime with a controlled number of emitters but with low reproducibility



# Acknowledgements

- PhD students & post-docs

Mickaël Busson, Laurent Lermusiaux, Nemanja Markesevic, Elise Gayet, Jeanne Heintz, Claudia Corti



- N. Bonod, J. Wenger & coworkers (Fresnel, Plasmonic resonators)
- G. Bellot & coworkers (CBS, DNA origamis)
- Sylvie Marguet (CEA, gold nanoparticles)

